

# *Research Article* **Dynamic Study of a Predator-Prey Model with Weak Allee Effect and Delay**

# **Yong Y[e](https://orcid.org/0000-0002-3377-1923) , <sup>1</sup> Hua Li[u](https://orcid.org/0000-0002-1135-4969) , <sup>1</sup> Yu-mei Wei,2 Ming Ma [,](https://orcid.org/0000-0003-1374-5346) <sup>1</sup> and Kai Zhan[g](https://orcid.org/0000-0002-7680-2849) <sup>1</sup>**

*1 School of Mathematics and Computer Science, Northwest Minzu University, Lanzhou 730000, China 2 Experimental Center, Northwest Minzu University, Lanzhou 730000, China*

Correspondence should be addressed to Hua Liu; 7783360@qq.com

Received 12 June 2019; Accepted 22 July 2019; Published 6 August 2019

Academic Editor: Ming Mei

Copyright © 2019 Yong Ye et al. Tis is an open access article distributed under the [Creative Commons Attribution License](https://creativecommons.org/licenses/by/4.0/), which permits unrestricted use, distribution, and reproduction in any medium, provided the original work is properly cited.

In this paper, a prey-predator model and weak Allee effect in prey growth and its dynamical behaviors are studied in detail. The existence, boundedness, and stability of the equilibria of the model are qualitatively discussed. Bifurcation analysis is also taken into account. Afer incorporating the searching delay and digestion delay, we establish a delayed predator-prey system with Allee effect. The results show that there exist stability switches and Hopf bifurcation occurs while the delay crosses a set of critical values. Finally, we present some numerical simulations to illustrate our theoretical analysis.

# **1. Introduction**

Some researchers have conducted extensive research on the dynamics of interacting prey-predator models to understand the long-term behavior of species. A wide variety of nonlinear coupled ordinary diferential equation models are proposed and analyzed for the interaction between prey and their predators. The classic predator-prey model is the Lotka-Volterra model, which was independently proposed by Lotka in the United States in 1925 and Volterra in Italy in 1926 [\[1](#page-13-0), [2\]](#page-13-1). The model was developed on the basis of a single-population growth model and has wide applicability. The mathematical form of the Lotka-Volterra model is

$$
\frac{dx}{dt} = rx - axy
$$
  

$$
\frac{dy}{dt} = cxy - my
$$
 (1)

In population dynamics,when the population density is very low, there is a positive correlation between the population unit growth rate and the population density. This phenomenon can be called the Allee efect [\[3](#page-13-2)[–5](#page-13-3)], starting with Allee's research [\[6](#page-13-4)]. The Allee effect is classified according to the density-dependent properties at low density. If the population density is low, a strong Allee efect will appear. If the proliferation rate is positive and increases, the Allee efect will be weak. Demographic Allee efects can be either weak or strong [\[7,](#page-13-5) [8\]](#page-13-6). When the density is below the critical threshold, the population afected by the strong Allee efect will have a negative average growth rate. Under deterministic dynamics, we fnd that populations that do not exceed this threshold will be extinct. Many jobs only consider the strong Allee efect, but in the work of Allee it is clear that the Allee efect also has a weak Allee effect [\[8](#page-13-6)-13].

Today, it is widely believed that the Allee effect greatly increases the likelihood of local and global extinction and can produce a rich variety of dynamic efects [\[14](#page-14-1)[–16](#page-14-2)]. And it is interesting and important to study the impact of Allee efect on the predator-prey models [\[17](#page-14-3)[–19](#page-14-4)]. In this paper, we introduced a predator-prey model with weak Allee effect:

<span id="page-0-0"></span>
$$
\frac{dx}{dt} = rx\left(1 - \frac{x}{K}\right)\frac{x}{x + A} - axy
$$
  

$$
\frac{dy}{dt} = c(ax)y - my
$$
  

$$
x(0) \ge 0
$$
  

$$
y(0) \ge 0
$$
 (2)

Here, the weak Allee effect term is  $P(x) = x/(x + A)$ , where  $A > 0$  is described as a "weak Allee effect constant" ([\[12\]](#page-14-5)).  $\hat{x}$  is the prey population and  $\hat{y}$  is the predator population,  $m$  is the intrinsic death rate of predators,  $c$  is the conversion efficiency from prey to predator,  $K$  is the carrying capacity,  $r$ is the intrinsic growth rate of prey, and  $a$  is the prey capture rate by their predators. It is more realistic to introduce time delay on the basis of traditional predator-prey model because it exists almost everywhere in biological activities and is considered as one of the reasons for the regular change of population density [\[20](#page-14-6)-26]. Therefore, in order to make the system established in this paper biologically closer to reality, incorporating the searching delay and digestion delay in the system [\(2\)](#page-0-0) is interesting. Based on the above considerations, We establish a predator-prey model with time delay and weak Allee efect, as follows:

$$
\frac{dx(t)}{dt} = rx(t)\left(1 - \frac{x(t)}{K}\right)\frac{x(t)}{x+A}
$$

$$
-ax(t-\tau_1)y(t-\tau_1)
$$

$$
\frac{dy(t)}{dt} = c(ax(t-\tau_2))y(t-\tau_2) - my(t)
$$

$$
x(0) \ge 0
$$

$$
y(0) \ge 0
$$
 (3)

where the time delay  $\tau_i$  (*i* = 1, 2) is the controlling or perturbed parameters,  $\tau_1$  is the searching delay, and  $\tau_2$  is the digestion delay.

The latter parts of the paper are described as follows. In Section [2,](#page-1-0) we discuss the boundedness, the stability of the equilibria, and bifurcation of the model [\(2\)](#page-0-0) in detail. In Section [3,](#page-4-0) we investigated local stability property of interior equilibrium point of the model [\(3\)](#page-1-1) with time delay; the Hopf bifurcation around the positive equilibrium point is also studied. In Section [4,](#page-8-0) we verify the previous theoretical derivation by numerical simulation.

# <span id="page-1-0"></span>**2. A Predator-Prey Model with Weak Allee Effect**

We easily see that model [\(2\)](#page-0-0) exhibits three equilibrium points  $E_0 = (0, 0), E_1 = (K, 0),$  and  $E_* = (x_*, y_*)$ . Here  $x_* =$ *m*/*ca*,  $y_* = (Krx_* - rx_*^2)/aK(x_* + A)$ . And for the positive equilibrium point(s), we have  $m / ca < K$ .

#### *2.1. Boundedness*

**Theorem 1.** *For the solution*  $(x(t), y(t))$  *of model,* 

$$
\limsup_{t \to \infty} \left( x(t) + \frac{1}{c} y(t) \right) \le \frac{K (m+r)^2}{4rm}.
$$
 (4)

*Proof.* We define  $\chi = x(t) + (1/c)y(t)$ . Then we can easily see that along the solution of system [\(2\),](#page-0-0)

$$
\frac{dx}{dt} = \frac{dx}{dt} + \frac{1}{c}\frac{dy}{dt}
$$
\n
$$
= rx\left(1 - \frac{x}{K}\right)\frac{x}{x + A} - axy + \frac{1}{c}c(ax)y - \frac{1}{c}my \quad (5)
$$
\n
$$
= rx\left(1 - \frac{x}{K}\right)\frac{x}{x + A} - \frac{m}{c}y.
$$

Thus, we see that for all large  $t > 0$ 

$$
\frac{dy}{dt} + m\chi = rx\left(1 - \frac{x}{K}\right)\frac{x}{x + A} - \frac{m}{c}y + mx + m\frac{1}{c}y
$$

$$
= rx\left(1 - \frac{x}{K}\right)\frac{x}{x + A} + mx
$$

$$
\le rx\left(1 - \frac{x}{K}\right) + mx = x\left(r + m - \frac{r}{K}x\right)
$$

$$
\le \frac{K}{4r}(m + r)^2.
$$

$$
(6)
$$

<span id="page-1-1"></span>Hence the standard comparison argument shows that

$$
\limsup_{t \to \infty} \left( x(t) + \frac{1}{c} y(t) \right) \le \frac{K (m+r)^2}{4rm}.
$$
 (7)

*2.2. Stability Analysis*

<span id="page-1-2"></span>**Theorem 2.** (1) Trivial equilibrium point  $E_0$  is always a saddle*node point.*

(2)  $E_1$  *is stable for a < m/cK and is a saddle point otherwise.*

*(3) Coexistence equilibrium*  $E_*$  *is locally asymptotically stable for*  $A < x_*^2/(K - 2x_*)$  *and is unstable node otherwise.* 

*Proof.* Let

$$
f(x, y) = rx\left(1 - \frac{x}{K}\right)\frac{x}{x + A} - axy
$$
  
g(x, y) = c(ax) y - my. (8)

So, the Jacobian matrix for the model [\(2\)](#page-0-0) is given by  $J =$  $\left(\begin{array}{c}\frac{\partial f}{\partial x} & \frac{\partial f}{\partial y} \\ \frac{\partial g}{\partial x} & \frac{\partial g}{\partial y}\end{array}\right)$  where

$$
\frac{\partial f}{\partial x} = \frac{-2rKx^3 + rKx^2 - 3rAx^2 + 2rAKx}{K(x + A)^2} - ay,
$$
  

$$
\frac{\partial f}{\partial y} = -ax,
$$
  

$$
\frac{\partial g}{\partial x} = cay,
$$
  

$$
\frac{\partial g}{\partial y} = cax - m.
$$
  
(9)

So we get

J  $= \left( \frac{-2rx^3 + rKx^2 - 3rAx^2 + 2rAKx}{K(x + A)^2} - ay - ax \right).$  (10)<br>cax - m).

First, it can be concluded by calculating the Jacobian matrix of the model [\(2\)](#page-0-0) at  $E_0$  given by

$$
J_0 = \begin{pmatrix} 0 & 0 \\ 0 & -m \end{pmatrix}.
$$
 (11)

And hence  $E_0$  is always a saddle-node point. Then, by evaluating the Jacobian matrix of the model [\(2\)](#page-0-0) at  $E_1$ , we find

$$
J_1 = \begin{pmatrix} \frac{-rK}{K+A} & -aK\\ 0 & caK-m \end{pmatrix}.
$$
 (12)

First eigenvalue  $-rK/(K + A)$  is negative; hence  $E_1$  is stable if  $caK − m < 0$  implying  $a < m/cK$ , and  $E_1$  is a saddle point when  $a > m/cK$ . Finally, the Jacobian matrix for the model [\(2\)](#page-0-0) evaluated at  $E_*$  is given by

$$
J_* = \left( \frac{-2rx_*^3 + rKx_*^2 - 3rAx_*^2 + 2rAKx_*}{K(x_* + A)^2} - ay_* - \frac{m}{c} \right). \tag{13}
$$

The characteristic polynomial is

$$
H(\lambda) = \lambda^2 - T\lambda + D \tag{14}
$$

where  $T = -rx_*(x_*^2 + 2Ax_* - AK)/K(x_*+A)^2$  and  $D = may_*$ . Thus, we have the following conclusions. (1) If  $T < 0$  and  $A < x_*^2/(K - 2x_*)$ , then the positive equilibrium is locally asymptotically stable. (2) If  $T > 0$  and  $A > x_*^2/(K - 2x_*)$ , then the positive equilibrium is unstable.

**Theorem 3.**  $E_1 = (K, 0)$  *is globally stable when*  $a < m/cK$ .

*Proof.* Consider the Lyapunov function:

$$
V(x, y) = \int_{K}^{x} \frac{u - K}{u} du + \frac{1}{c} y.
$$
 (15)

The derivative of  $V$  along the solution of the model is

$$
\dot{V} = \frac{x - K}{x} \frac{dx}{dt} + \frac{1}{c} \frac{dy}{dt}
$$
\n
$$
= \frac{x - K}{x} \left[ rx \left( 1 - \frac{x}{K} \right) \frac{x}{x + A} - axy \right]
$$
\n
$$
+ \frac{1}{c} \left[ c \left( ax \right) y - my \right]
$$
\n
$$
= (x - K) \left[ r \left( 1 - \frac{x}{K} \right) \frac{x}{x + A} \right] - a \left( x - K \right) y + axy \quad (16)
$$
\n
$$
- \frac{m}{c} y
$$
\n
$$
= (x - K) \frac{rKx - rx^2}{K(x + A)} - a \left( x - K \right) y + axy - \frac{m}{c} y
$$
\n
$$
= \frac{-rx(x - K)^2}{K(x + A)} + aKy - \frac{m}{c} y \le aKy - \frac{m}{c} y.
$$

#### *2.3. Bifurcation Analysis*

## *2.3.1. Transcritical Bifurcation*

**Theorem 4.** *The model enters into transcritical bifurcation around*  $E_1$  *at*  $a = a_0$ *, where*  $a_0 = m/cK$ *.* 

*Proof.* One of the eigenvalues of  $J_1$  will be zero if  $J_1 = 0$  which gives  $a = a_0$ . At this point, the other eigenvalue is  $-rK/(K +$ A). If  $V$  and  $W$  denote the eigenvectors corresponding to the eigenvalue 0 of the matrices  $I_1$  and  $I_1^T$ , respectively, then we obtain  $V = (-a(K + A)/r, 1)^T$  and  $W = (0, 1)^T$ , where  $J_1^T =$  $\binom{-rK/(K+A)}{-aK}$  0,  $V_1 = -a(K+A)/r$ ,  $V_2 = 1$ .

$$
W^{T} f_{a} (\overline{x}, \overline{y}, a_{0}) = 0,
$$
  
\n
$$
W^{T} [Df_{a} (\overline{x}, \overline{y}, a_{0}) V] = cK \neq 0,
$$
  
\n
$$
W^{T} [D^{2} f (\overline{x}, \overline{y}, a_{0}) (V, V)]
$$
  
\n
$$
= W^{T} \left( \frac{\partial^{2} f_{1}}{\partial x^{2}} V_{1}^{2} + 2 \frac{\partial^{2} f_{1}}{\partial x \partial y} V_{1} V_{2} + \frac{\partial^{2} f_{1}}{\partial y^{2}} V_{2}^{2} \right)
$$
  
\n
$$
= W^{T} \left( \frac{\partial^{2} f_{2}}{\partial x^{2}} V_{1}^{2} + 2 \frac{\partial^{2} f_{2}}{\partial x \partial y} V_{1} V_{2} + \frac{\partial^{2} f_{2}}{\partial y^{2}} V_{2}^{2} \right)_{(\overline{x}, \overline{y}, a_{0})}
$$
  
\n
$$
= \frac{-2a^{2} (K + A)}{r} \neq 0.
$$
 (17)

Therefore, by the Sotomayor theorem, we can find that the model experiences transcritical bifurcation at  $a = a_0$  around the axial equilibrium  $E_1$ . the axial equilibrium  $E_1$ .

2.3.2. Hopf Bifurcation. From Theorem [2,](#page-1-2) model [\(2\)](#page-0-0) undergoes bifurcation if  $A = x_*^2/(K - 2x_*)$ . The purpose of this section is to show that model [\(2\)](#page-0-0) undergoes a Hopf bifurcation if  $A = x_*^2/(K - 2x_*)$ . We analyze the Hopf bifurcation occurring at  $E_* = (x_*, y_*)$  by choosing as the bifurcation parameter. Denote

$$
A_0 = \frac{{x_*}^2}{K - 2x_*}.
$$
 (18)

When  $A = A_0$ , we have  $T = -rx_*(x^2 + 2Ax^2 - AK)/K(x^2 +$  $(A)^2 = 0$ . Thus, the Jacobian matrix  $J_*$  has a pair of imaginary eigenvalues  $\lambda = \pm i \sqrt{m a y_*}$ . Let  $\lambda = \alpha(A) \pm \beta(A) i$  be the roots of  $\lambda^2 - T\lambda + D = 0$ ; then

$$
\alpha^{2} - \beta^{2} - \alpha T + D = 0
$$
  
2\alpha\beta - T\beta = 0 (19)

and

$$
\alpha = \frac{T}{2}
$$
\n
$$
\beta = \frac{\sqrt{4D - T^2}}{2}
$$
\n
$$
\Big|_{A = A_0} = \frac{-rx_*^3}{2AK(x_* + A)^2} < 0
$$
\n(20)

By the Poincare-Andronov-Hopf Bifurcation Theorem, we know that model [\(2\)](#page-0-0) undergoes a Hopf bifurcation at  $E_*$  =  $(x_*, y_*)$  when  $A = A_0$ . However, the detailed nature of the Hopf Bifurcation needs further analysis of the normal form of the model. Set  $x = X + x_*$  and  $y = Y + y_*$ , to  $(x_*, y_*)$  as origin of coordinates  $(X, Y)$ . We have the following model:

 $d\alpha$  $\overline{dA}$ 

$$
\frac{dX}{dt} = a_{11}X + a_{12}Y + F_1(X, Y)
$$
\n
$$
\frac{dY}{dt} = a_{21}X + a_{22}Y + F_2(X, Y)
$$
\n(21)

where  $a_{11} = (-2rx_*^3 + rKx_*^2 - 3rAx_*^2 + 2rAKx_*)/K(x_* +$  $(A)^{2} - ay_{*}, a_{12} = -m/c, a_{21} = cay_{*}, a_{22} = 0$ , and

$$
F_1(X, Y)
$$
  
=  $A_1 X^2 + A_2 XY + A_3 Y^2 + B_1 X^3 + B_2 X^2 Y + B_3 XY^2$   
+  $B_4 Y^3 + P_1(X, Y)$   

$$
F_2(X, Y)
$$
  
=  $C_1 X^2 + C_2 XY + C_3 Y^2 + D_1 X^3 + D_2 X^2 Y + D_3 XY^2$   
+  $D_4 Y^3 + P_2(X, Y)$ 

$$
A_{1} = \frac{-rx_{*}^{3} - 3rAx_{*}^{2} - 3rA^{2}x_{*} + rA^{2}K}{K(x_{*} + A)^{3}},
$$
\n
$$
A_{2} = -\frac{a}{2},
$$
\n
$$
A_{3} = 0
$$
\n
$$
B_{1}
$$
\n
$$
= \frac{rx_{*}^{3} + (3rA - r)x_{*}^{2} + (3rA^{2} - 2rA)x_{*} - (rA^{2} + rA^{2}K)}{2K(x_{*} + A)^{3}},
$$
\n
$$
B_{2} = 0,
$$
\n
$$
B_{3} = 0,
$$
\n
$$
B_{4} = 0
$$
\n
$$
C_{1} = 0,
$$
\n
$$
C_{2} = \frac{ca}{2},
$$
\n
$$
C_{3} = 0
$$
\n
$$
D_{1} = 0,
$$
\n
$$
D_{2} = 0,
$$
\n
$$
D_{3} = 0,
$$
\n
$$
D_{4} = 0
$$
\n(22)\n
$$
D_{5} = 0,
$$
\n
$$
D_{6} = 0
$$
\n(22)

where  $P_1(X, Y)$  and  $P_2(X, Y)$  are smooth functions of X and Y at least of order four. Now, using the transformation  $u=X$ ,  $v = -(1/\beta)(a_{11}X + a_{12}Y)$ , we obtain

$$
\frac{du}{dt} = -\beta v + G_1 (u, v)
$$
  

$$
\frac{dv}{dt} = \beta u + G_2 (u, v)
$$
 (23)

where

$$
G_{1}(u, v) = F_{1}\left(u, -\frac{1}{a_{12}}(a_{11}u + \beta v)\right)
$$
  
\n
$$
G_{2}(u, v) = -\frac{1}{\beta}\left(a_{11}F_{1}\left(u, -\frac{1}{a_{12}}(a_{11}u + \beta v)\right) - (24) + a_{12}F_{2}\left(u, -\frac{1}{a_{12}}(a_{11}u + \beta v)\right)\right)
$$

so

$$
G_{1}(u, v) = A_{1}u^{2} + A_{2}u\left(-\frac{1}{a_{12}}(a_{11}u + \beta v)\right) + B_{1}u^{3}
$$
  
\n
$$
G_{2}(u, v) = -\frac{1}{\beta}\left[a_{11}\left(A_{1}u^{2}\right) + A_{2}u\left(-\frac{1}{a_{12}}(a_{11}u + \beta v)\right) + B_{1}u^{3}\right)\right]
$$
\n
$$
+ a_{12}uC_{2}\left(-\frac{1}{a_{12}}(a_{11}u + \beta v)\right)
$$
\n(25)

Set

$$
\sigma = \frac{1}{16} \left[ \frac{\partial^3 G_1}{\partial u^3} + \frac{\partial^3 G_1}{\partial u \partial v^2} + \frac{\partial^3 G_2}{\partial u^2 \partial v} + \frac{\partial^3 G_2}{\partial v^3} \right] \n+ \frac{1}{16\beta} \left[ \frac{\partial^2 G_1}{\partial u \partial v} \left( \frac{\partial^2 G_1}{\partial u^2} + \frac{\partial^2 G_1}{\partial v^2} \right) - \frac{\partial^2 G_2}{\partial u \partial v} \left( \frac{\partial^2 G_2}{\partial u^2} + \frac{\partial^2 G_2}{\partial v^2} \right) - \frac{\partial^2 G_1}{\partial u^2} \frac{\partial^2 G_2}{\partial u^2} + \frac{\partial^2 G_1}{\partial v^2} \frac{\partial^2 G_2}{\partial v^2} \right]
$$
\n(26)

where

$$
\frac{\partial^3 G_1}{\partial u^3} = 6B_1,
$$
\n
$$
\frac{\partial^3 G_1}{\partial u \partial v^2} = 0,
$$
\n
$$
\frac{\partial^3 G_2}{\partial u^2 \partial v} = 0,
$$
\n
$$
\frac{\partial^3 G_2}{\partial v^3} = 0,
$$
\n
$$
\frac{\partial^2 G_1}{\partial u \partial v} = \frac{A_2 \beta}{a_{12}},
$$
\n
$$
\frac{\partial^2 G_2}{\partial u \partial v} = \frac{A_2 a_{11}}{a_{12}},
$$
\n
$$
\frac{\partial^2 G_1}{\partial v^2} = 0,
$$
\n
$$
\frac{\partial^2 G_2}{\partial v^2} = 0,
$$
\n
$$
\frac{\partial^2 G_2}{\partial v^2} = 2\left(A_1 - \frac{A_2 a_{11}}{a_{12}}\right) + 6B_1 u,
$$
\n
$$
\frac{\partial^2 G_2}{\partial u^2} = 2\left(\frac{A_2 a_{11}^2}{\beta a_{12}} - \frac{A_1 a_{11}}{\beta}\right) - \frac{6B_1 a_{11}}{\beta} u.
$$

So

$$
\sigma = \frac{3B_1}{8} + \frac{1}{16\beta} \left( \frac{A_2\beta}{a_{12}} \left( 2\left(A_1 - \frac{A_2a_{11}}{a_{12}}\right) + 6B_1u \right) - \frac{A_2a_{11}}{a_{12}} \left( 2\left(\frac{A_2a_{11}^2}{\beta a_{12}} - \frac{A_1a_{11}}{\beta}\right) - \frac{6B_1a_{11}}{\beta}u \right) - \left( 2\left(A_1 - \frac{A_2a_{11}}{a_{12}}\right) + 6B_1u \right) - \left( 2\left(\frac{A_2a_{11}^2}{\beta a_{12}} - \frac{A_1a_{11}}{\beta}\right) - \frac{6B_1a_{11}}{\beta}u \right) \right)
$$
\n(28)

If  $\sigma$  < 0, the equilibrium  $E_*$  is destabilized through a Hopf bifurcation that is supercritical and Hopf bifurcation is subcritical otherwise.

# <span id="page-4-0"></span>**3. Delayed Model with Weak Allee Effect**

Let  $X(t) = x(t) - x_*, Y(t) = y(t) - y_*$ ; then the model [\(3\)](#page-1-1) can be expressed as in the following matrix form afer linearization:

$$
\frac{d}{dt} \begin{pmatrix} X(t) \\ Y(t) \end{pmatrix}
$$
\n
$$
= A_1 \begin{pmatrix} X(t) \\ Y(t) \end{pmatrix} + A_2 \begin{pmatrix} X(t - \tau_1) \\ Y(t - \tau_1) \end{pmatrix}
$$
\n
$$
+ A_3 \begin{pmatrix} X(t - \tau_2) \\ Y(t - \tau_2) \end{pmatrix},
$$
\n
$$
A_1 \tag{29}
$$
\n
$$
\begin{pmatrix} -2rx_*^3 + rKx_*^2 - 3rAx_*^2 + 2rAKx_* \\ 0 \end{pmatrix}
$$

$$
= \begin{pmatrix} \frac{-2rx_*^3 + rKx_*^2 - 3rAx_*^2 + 2rAKx_*}{K(x_* + A)^2} & 0\\ 0 & -m \end{pmatrix},
$$
  
\n
$$
A_2 = \begin{pmatrix} -ay_* & -ax_*\\ 0 & 0 \end{pmatrix},
$$
  
\n
$$
A_3 = \begin{pmatrix} 0 & 0\\ cap_* & cax_* \end{pmatrix}.
$$

3.1. Stability Analysis. The characteristic polynomial is

$$
H(\lambda) = \lambda^2 - T\lambda + D \tag{30}
$$

where

$$
T = \frac{-2rx_*^3 + rKx_*^2 - 3rAx_*^2 + 2rAKx_*}{K(x_* + A)^2}
$$
  
\n
$$
-ay_*e^{-\lambda\tau_1} + cav_*e^{-\lambda\tau_2} - m,
$$
  
\n
$$
D = \left(\frac{-2rx_*^3 + rKx_*^2 - 3rAx_*^2 + 2rAKx_*}{K(x_* + A)^2}\right)
$$
  
\n
$$
-ay_*e^{-\lambda\tau_1}\left((cx_*e^{-\lambda\tau_2} - m) + ca^2x_*y_*e^{-\lambda\tau_1}e^{-\lambda\tau_2}\right).
$$
  
\n(31)

Let

$$
G = \frac{\left(-2rx_*^3 + rKx_*^2 - 3rAx_*^2 + 2rAKx_*\right)}{K\left(x_* + A\right)^2} \tag{32}
$$

So

$$
H(\lambda) = \lambda^2 - \left(G - ay_*e^{-\lambda \tau_1} + ca x_* e^{-\lambda \tau_2} - m\right)\lambda
$$
  
- 
$$
mG + may_*e^{-\lambda \tau_1} + ca x_* e^{-\lambda \tau_2}G
$$
  
Let  $\lambda(\tau) = \eta(\tau) + i\omega(\tau), \eta(\tau_0) = 0$ , and  $\omega(\tau_0) = \omega_0$ . (33)

**Theorem 5.** Assume  $A < x_*^2/(K - 2x_*)$ , when  $\tau_1 > 0$ ,  $\tau_2 = 0$ ; *we have the following conclusions. (1) When*  $4FG - 2Fm + G^2 +$ 

 $F^2 + m^2 - D^2 > 0$  and  $-2FmG^2 + F^2G^2 + m^2G^2 - E^2 > 0$ , *the positive equilibrium*  $E_* = (x_*, y_*)$  *is locally asymptotically stable.* (2) Hopf bifurcation occurs when  $\tau_1$  passes the critical *value*

$$
\overline{\tau}_{1} = \frac{1}{\omega_{0}}
$$
  
arccos 
$$
\frac{EGm - EFG - \omega_{0}^{2}mD + \omega_{0}^{2}FD + \omega_{0}^{2}GD + \omega_{0}^{2}E}{\omega_{0}^{2}D^{2} + E^{2}}.
$$
 (34)

*Proof.* The characteristic equation is

$$
\lambda^{2} - \left(G - ay_{*}e^{-\lambda \tau_{1}} + ca x_{*} - m\right)\lambda - mG
$$
  
+
$$
may_{*}e^{-\lambda \tau_{1}} + ca x_{*}G = 0.
$$
 (35)

Next we suppose that  $\lambda(\tau_0) = i\omega_0$  is a solution of  $H(\lambda)$  for some  $\tau > 0$ ; then we have

$$
- \omega_0^2 - i\omega_0 \left( G - a y_* e^{-i\omega_0 \tau_1} + c a x_* - m \right) - mG
$$
  
+ 
$$
m a y_* e^{-i\omega_0 \tau_1} + c a x_* G = 0.
$$
 (36)

Then

$$
- \omega_0^2 - i\omega_0 (G - De^{-i\omega_0 \tau_1} + F - m) - mG + E e^{-i\omega_0 \tau_1}
$$
  
+ FG = 0. (37)

Where  $D = ay_*$ ,  $E = ma y_*$ ,  $F = ca x_*$  we know

$$
e^{-i\omega_0 \tau} = \cos \omega_0 \tau - i \sin \omega_0 \tau.
$$
 (38)

So we get

$$
-\omega_0^2 - i\omega_0 G + i\omega_0 D \cos \omega_0 \tau_1 + \omega_0 D \sin \omega_0 \tau_1 - i\omega_0 F
$$

$$
+ i\omega_0 m - mG + E \cos \omega_0 \tau_1 - iE \sin \omega_0 \tau_1 + FG \qquad (39)
$$

$$
= 0.
$$

Separate real and imaginary parts

$$
-\omega_0^2 + \omega_0 D \sin \omega_0 \tau_1 - mG + E \cos \omega_0 \tau_1 + FG = 0
$$
  

$$
-\omega_0 G + \omega_0 D \cos \omega_0 \tau_1 - \omega_0 F + \omega_0 m - E \sin \omega_0 \tau_1
$$
 (40)  

$$
= 0
$$

Then

$$
\omega_0 D \sin \omega_0 \tau_1 + E \cos \omega_0 \tau_1 = {\omega_0}^2 - FG + mG
$$
  

$$
\omega_0 D \cos \omega_0 \tau_1 - E \sin \omega_0 \tau_1 = \omega_0 G + \omega_0 F - \omega_0 m
$$
 (41)

So

$$
\omega_0^4 + \left(4FG - 2Fm + G^2 + F^2 + m^2 - D^2\right)\omega_0^2
$$
  
- 2FmG<sup>2</sup> + F<sup>2</sup>G<sup>2</sup> + m<sup>2</sup>G<sup>2</sup> - E<sup>2</sup> = 0 (42)

We assume that

$$
W = \omega_0^2
$$
  
\n
$$
W^2 + (4FG - 2Fm + G^2 + F^2 + m^2 - D^2)W
$$
 (43)  
\n
$$
-2FmG^2 + F^2G^2 + m^2G^2 - E^2 = 0.
$$

If  $4FG - 2Fm + G^2 + F^2 + m^2 - D^2 > 0$  and  $-2FmG^2 + F^2G^2 +$  $m^2G^2 - E^2 > 0$ , then all roots of equation have negative real parts for all  $\tau_1 > 0$ ,  $\tau_2 = 0$ ; that is, the equilibrium  $E_*$  $(x_*, y_*)$  is locally asymptotically stable:

W

$$
=\frac{-M\pm\sqrt{M^2-4\left(-2FmG^2+F^2G^2+m^2G^2-E^2\right)}}{2}\tag{44}
$$

where

$$
M = 4FG - 2Fm + G2 + F2 + m2 - D2.
$$
 (45)

If  $-M > 0$ ,  $M^2 = 4(2FmG^2 + F^2G^2 + m^2G^2 - E^2)$  and  $-2FmG^2 +$  $F^2G^2 + m^2G^2 - E^2 > 0$ , there is a unique positive solution; the equilibrium  $E_* = (x_*, y_*)$  is unstable. Also if  $-M > 0$ ,  $M^2 > 4(-2FmG^2 + F^2G^2 + m^2G^2 - E^2)$ , and  $-2FmG^2 + F^2G^2 +$  $m^2G^2 - E^2 > 0$ , then there are two positive solutions. We have

$$
-\omega_0^2 + \omega_0 D \sin \omega_0 \tau_1 - mG + E \cos \omega_0 \tau_1 + FG = 0 \quad (46)
$$

So

$$
\sin \omega_0 \tau_1 = \frac{\omega_0^2 + mG - E \cos \omega_0 \tau_1 - FG}{\omega_0 D} \tag{47}
$$

Then, we get

 $\cos \omega_0 \tau_1$ 

$$
= \frac{EGm - EFG - \omega_0^2 mD + {\omega_0}^2 FD + {\omega_0}^2 GD + {\omega_0}^2 E}{\omega_0^2 D^2 + E^2}.
$$
 (48)

It shows that if  $-M > 0$  and  $-2FmG^2 + F^2G^2 + m^2G^2 - E^2 > 0$ ,

$$
\lambda^{2} - \left(G - ay_{*}e^{-\lambda \tau_{1}} + ca x_{*} + m\right)\lambda + mG
$$
  
-
$$
may_{*}e^{-\lambda \tau_{1}} + ca x_{*}G = 0
$$
 (49)

has a pair of imaginary eigenvalues  $\lambda=\pm i\omega_0, \eta(\overline{\tau}_1)=0$ when  $\overline{\tau}_{1j}^{\ \pm} = (1/\omega_0^{\ \pm}) \arccos((EGm - EFG - \omega_0^2mD +$  $\omega_0^2 FD + \omega_0^2 GD + \omega_0^2 E / (\omega_0^2 D^2 + E^2) + 2\pi j / \omega_0^{\pm}, j =$  $0, 1, 2, \cdots$ .

Next verify the cross-sectional conditions:

$$
\left(\frac{d\text{Re}\left(\lambda\right)}{d\tau_{1}}\right)^{-1} \neq 0\tag{50}
$$

According to

$$
\lambda^{2} - \left(G - ay_{*}e^{-\lambda \tau_{1}} + cax_{*} - m\right)\lambda - mG + may_{*}e^{-\lambda \tau_{1}}+ cax_{*}G = 0
$$
\n(51)

At this time,  $\tau_1 = \overline{\tau}_1$ , where  $\overline{\tau}_1$  is the value of  $\overline{\tau}_{1i}$  at  $j = 0$ . We get

$$
2\lambda \frac{d\lambda}{d\tau_1} - \frac{d\lambda}{d\tau_1} \left( G - De^{-\lambda \tau_1} + F - m \right)
$$
  
+ 
$$
\lambda De^{-\lambda \tau_1} \left( -\tau_1 \frac{d\lambda}{d\tau_1} - \lambda \right)
$$
  
+ 
$$
E e^{-\lambda \tau_1} \left( -\tau_1 \frac{d\lambda}{d\tau_1} - \lambda \right) = 0
$$
 (52)

Then  
\n
$$
\left(\frac{d\lambda}{d\tau_1}\right)^{-1}
$$
\n
$$
= -\frac{\tau_1}{\lambda} + \frac{D}{\lambda (E + \lambda D)}
$$
\n
$$
+ \frac{[2\lambda - (G + F - m)]}{\lambda [\lambda^2 - \lambda (G + F - m) - mG + FG]}.
$$
\n
$$
\left(\frac{dRe(\lambda)}{d\tau_1}\right)^{-1}
$$
\n
$$
= -\frac{D^2}{\omega_0^2 D^2 + E^2}
$$
\n
$$
+ \frac{-(G + F - m)^2 + 2(-mG + FG - \omega_0^2)}{(G + F - m)^2 \omega_0^2 + (-mG + FG - \omega_0^2)^2}
$$
\n
$$
\neq 0.
$$

**Theorem 6.** Assume  $A < x_*^2/(K - 2x_*)$ , when  $\tau_1 = 0, \tau_2 > 0$ ; *we have the following conclusions.* (1) If  $-2EmG - F^2G^2$  +  $m^2G^2 + E^2 > 0$ , the positive equilibrium  $E_* = (x_*, y_*)$  is locally asymptotically stable. (2) When –2EmG– $F^2G^2$ +m $^2G^2$ +E $^2$  < 0, if  $\tau_2 < \overline{\tau}_2$ , the positive equilibrium  $E_* = (x_*, y_*)$  is locally *asymptotically stable; if*  $\tau_2 > \bar{\tau}_2$ *, the positive equilibrium is unstable.* (3) Hopf bifurcation occurs when  $\tau$ <sub>2</sub> passes the critical *value*

$$
\overline{\tau}_2 = \frac{1}{\omega_0} \arccos \frac{{\omega_0}^2 (-m - D) F + FG (E - mG)}{{-\omega_0}^2 F^2 - E^2}.
$$
 (54)

Proof. The characteristic equation is

$$
\lambda^{2} - (G - ay_{*} + ca x_{*} e^{-\lambda \tau_{2}} + m) \lambda + mG - may_{*}
$$
  
+ 
$$
ca x_{*} e^{-\lambda \tau_{2}} G = 0
$$
 (55)

Next we suppose that  $\lambda(\tau_0) = i\omega_0$  is a solution of  $H(\lambda)$  for some  $\tau > 0$ ; then we have

$$
-\omega_0^2 - i\omega_0 (G - ay_* + cax_*e^{-i\omega_0 \tau_2} - m) - mG
$$
  
+
$$
may_* + cax_*Ge^{-i\omega_0 \tau_2} = 0.
$$
 (56)

Then

$$
-\omega_0^2 - i\omega_0 (G - D + Fe^{-i\omega_0 \tau_2} - m) - mG + E
$$
  
+  $FGe^{-i\omega_0 \tau_2} = 0.$  (57)

Where  $D = ay_*, E = may_*, F = ca x_*$  we know

$$
e^{-i\omega_0 \tau} = \cos \omega_0 \tau - i \sin \omega_0 \tau.
$$
 (58)

So we get

$$
-\omega_0^2 - i\omega_0 G + i\omega_0 D - \omega_0 F \sin \omega_0 \tau_2 - i\omega_0 F \cos \omega_0 \tau_2
$$
  
+  $i\omega_0 m - mG + E - iFG \sin \omega_0 \tau_2$  (59)  
+  $FG \cos \omega_0 \tau_2 = 0$ .

Separate real and imaginary parts

$$
-\omega_0^2 - \omega_0 F \sin \omega_0 \tau_2 - mG + E + FG \cos \omega_0 \tau_2 = 0
$$
  

$$
-\omega_0 G + \omega_0 D - \omega_0 F \cos \omega_0 \tau_2 + \omega_0 m - FG \sin \omega_0 \tau_2
$$
 (60)  

$$
= 0
$$

Then

$$
-\omega_0 F \sin \omega_0 \tau_2 + FG \cos \omega_0 \tau_2 = {\omega_0}^2 - E + mG
$$
  

$$
-\omega_0 F \cos \omega_0 \tau_2 - FG \sin \omega_0 \tau_2 = \omega_0 G - \omega_0 D - \omega_0 m
$$
 (61)

So

$$
\omega_0^{-4}
$$

+
$$
(-2E - 2DG + 2Dm + G^2 - F^2 + m^2 + D^2) \omega_0^2
$$
 (62)  
- $2EmG - F^2G^2 + m^2G^2 + E^2 = 0$ 

We assume that

$$
W = \omega_0^2
$$
  
W<sup>2</sup> + (-2E - 2DG + 2Dm + G<sup>2</sup> - F<sup>2</sup> + m<sup>2</sup> + D<sup>2</sup>)W (63)  
- 2EmG - F<sup>2</sup>G<sup>2</sup> + m<sup>2</sup>G<sup>2</sup> + E<sup>2</sup> = 0

We can easily find

$$
-2E - 2DG + 2Dm + G^2 - F^2 + m^2 + D^2 > 0.
$$
 (64)

If  $-2EmG - F^2G^2 + m^2G^2 + E^2 > 0$ , the positive equilibrium  $E_* = (x_*, y_*)$  is locally asymptotically stable,−2EmG– $F^2G^2 +$  $m^2G^2 + E^2 < 0$ , and the positive equilibrium is unstable. We have

$$
-\omega_0 F \sin \omega_0 \tau_2 + FG \cos \omega_0 \tau_2 = {\omega_0}^2 - E + mG \tag{65}
$$

So

$$
\sin \omega_0 \tau_1 = \frac{\omega_0^2 - E + mG - FG \cos \omega_0 \tau_2}{-\omega_0 F} \tag{66}
$$

Then, we get

$$
\cos \omega_0 \tau_2 = \frac{\omega_0^2 (-m - D) F - FG (-E + mG)}{-\omega_0^2 F^2 - E^2}.
$$
 (67)

It shows that if–2*EmG* –  $F^2G^2 + m^2G^2 + E^2 < 0$ ,

$$
\lambda^{2} - \left(G - ay_{*} + ca x_{*} e^{-\lambda \tau_{2}} - m\right) \lambda - mG + m a y_{*}
$$

$$
+ ca x_{*} e^{-\lambda \tau_{2}} G = 0
$$
\n(68)

has a pair of imaginary eigenvalues  $\lambda = \pm i\omega_0$ ,  $\eta(\overline{\tau}_2)=0$ when  $\overline{\tau}_{2j}^{\pm} = (1/\omega_0^{\pm}) \arccos((\omega_0^2(-m - D)F + FG(E (mG)) / (-\omega_0^2 F^2 - E^2) + 2\pi j / \omega_0^{\pm}, \ \ j = 0, 1, 2, \cdots.$ Next verify the cross-sectional conditions:

$$
\left(\frac{d\text{Re}\left(\lambda\right)}{d\tau_{2}}\right)^{-1} \neq 0\tag{69}
$$

According to

$$
\lambda^{2} - (G - ay_{*} + ca x_{*} e^{-\lambda \tau_{2}} - m) \lambda - mG + may_{*}
$$
  
+ 
$$
ca x_{*} e^{-\lambda \tau_{2}} G = 0
$$
 (70)

At this time,  $\tau_2 = \overline{\tau}_2$ , where  $\overline{\tau}_2$  is the value of  $\overline{\tau}_2$  at  $j = 0$ . We get

$$
2\lambda \frac{d\lambda}{d\tau_2} - \frac{d\lambda}{d\tau_2} \left( G - D + Fe^{-\lambda \tau_2} - m \right)
$$

$$
- \lambda Fe^{-\lambda \tau_2} \left( -\tau_2 \frac{d\lambda}{d\tau_2} - \lambda \right)
$$

$$
+ FG e^{-\lambda \tau_2} \left( -\tau_2 \frac{d\lambda}{d\tau_2} - \lambda \right) = 0
$$

$$
(71)
$$

Then

$$
\left(\frac{d\lambda}{d\tau_2}\right)^{-1} \n= -\frac{\tau_2}{\lambda} - \frac{F}{\lambda (G - \lambda)} \n+ \frac{2\lambda - (G - D - m)}{\lambda [\lambda^2 - \lambda (G - D - m) - mG + E]} \n\left(\frac{dRe(\lambda)}{d\tau_2}\right)^{-1} \n= -\frac{F}{\omega_0^2 + G^2} \n+ \frac{2\left[(-mG + E) - \omega_0^2\right] - (G - D - m)^2}{\omega_0^2 (G - D - m)^2 + \left[(-mG + E) - \omega_0^2\right]^2} \n\neq 0.
$$

**Theorem 7.** Assume  $A < x_*^2/(K - 2x_*)$ , when  $\tau_{1, \tau} \neq 0, \tau_2 \neq 0$ ; *we have the following conclusions. (1) When and*  $G^2 - F^2 + m^2 D^2 + 2DF > 0$  and  $-2EFG - F^2G^2 + m^2G^2 - E^2 > 0$ , the positive *equilibrium*  $E_* = (x_*, y_*)$  *is locally asymptotically stable.* (2) *Hopf bifurcation occurs when*  $\tau_1$  *passes the critical value* 

$$
\overline{\tau} = \frac{1}{\omega_0}
$$
\n
$$
\text{arctan}\frac{\left(\omega_0^2 + m\overline{G}\right)\left(\omega_0 D - \omega_0 F\right) - \omega_0 \left(G - m\right)}{-\left(\omega_0^2 + m\overline{G}\right)\left(FG + E\right) + \omega_0 \left(G - m\right)}.
$$
\n(73)

*Proof.* The characteristic equation is

$$
\lambda^{2} - (G - ay_{*}e^{-\lambda \tau_{1}} + cax_{*}e^{-\lambda \tau_{2}} - m)\lambda - mG
$$
  
+
$$
may_{*}e^{-\lambda \tau_{1}} + cax_{*}e^{-\lambda \tau_{2}}G = 0.
$$
 (74)

Next we suppose that  $\lambda(\tau_0) = i\omega_0$  is a solution of  $H(\lambda)$  for some  $\tau > 0$ ; then we have

$$
-\omega_0^2 - i\omega_0 \left(G - De^{-i\omega_0 \tau_1} + Fe^{-i\omega_0 \tau_2} + m\right) - mG
$$
  
+  $Be^{-i\omega_0 \tau_1} + FGe^{-i\omega_0 \tau_2} = 0.$  (75)

Where  $D = ay_*, E = may_*, F = ca x_*$  we know

$$
e^{-i\omega_0 \tau} = \cos \omega_0 \tau - i \sin \omega_0 \tau.
$$
 (76)

So we get

$$
-\omega_0^2 - i\omega_0 G + i\omega_0 D \cos \omega_0 \tau_1 + \omega_0 D \sin \omega_0 \tau_1
$$
  

$$
-\omega_0 F \sin \omega_0 \tau_2 - i\omega_0 F \cos \omega_0 \tau_2 + i\omega_0 m - mG
$$
  
+ E \cos \omega\_0 \tau\_1 - iE \sin \omega\_0 \tau\_1 - iFG \sin \omega\_0 \tau\_2  
+ FG \cos \omega\_0 \tau\_2 = 0 (77)

Separate real and imaginary parts

Let  $\tau_1 = \tau_2 = \tau > 0$ . Then

$$
-\omega_0^2 + \omega_0 D \sin \omega_0 \tau_1 - \omega_0 F \sin \omega_0 \tau_2 - mG
$$
  
+ E cos  $\omega_0 \tau_1$  + FG cos  $\omega_0 \tau_2$  = 0  

$$
-\omega_0 G + \omega_0 D \cos \omega_0 \tau_1 - \omega_0 F \cos \omega_0 \tau_2 + \omega_0 m
$$

$$
- FG \sin \omega_0 \tau_2 - E \sin \omega_0 \tau_1 = 0
$$
(78)

$$
\omega_0 D \sin \omega_0 \tau - \omega_0 F \sin \omega_0 \tau + E \cos \omega_0 \tau
$$
  
+  $FG \cos \omega_0 \tau = \omega_0^2 + mG$   

$$
\omega_0 D \cos \omega_0 \tau - \omega_0 F \cos \omega_0 \tau - FG \sin \omega_0 \tau
$$
  
-  $E \sin \omega_0 \tau = \omega_0 G - \omega_0 m$  (79)

So

$$
\omega_0^4 + \left(G^2 - F^2 + m^2 - D^2 + 2DF\right)\omega_0^2 - 2EFH
$$
  

$$
-F^2G^2 + m^2G^2 - E^2 = 0
$$
 (80)

We assume that

$$
W = \omega_0^2
$$
  
\n
$$
W^2 + (G^2 - F^2 + m^2 - D^2 + 2DF)W - 2EFH
$$
 (81)  
\n
$$
-F^2G^2 + m^2G^2 - E^2 = 0
$$

If  $G^2 - F^2 + m^2 - D^2 + 2DF > 0$  and  $-2EFH - F^2G^2 + m^2G^2 - E^2 >$ 0, then all roots of equation have negative real parts for all  $\tau_1 > 0$ ,  $\tau_2 > 0$ ; that is, the equilibrium  $E_* = (x_*, y_*)$  is locally asymptotically stable:

$$
W = \frac{-N \pm \sqrt{N^2 - 4\left(-2EFH - F^2G^2 + m^2G^2 - E^2\right)}}{2} \tag{82}
$$

where

$$
N = G2 - F2 + m2 - D2 + 2DF
$$
 (83)

If  $-N > 0$ ,  $N^2 = 4(-2EFH - F^2G^2 + m^2G^2 - E^2)$ , and  $-2EFH F^2G^2 + m^2G^2 - E^2 > 0$ , there is a unique positive solution; the equilibrium  $E_* = (x_*, y_*)$  is unstable. Also if  $-N > 0$ ,  $N^2 > 4(-2EFH - F^2G^2 + m^2G^2 - E^2)$ , and  $-2EFH - F^2G^2 +$  $m^2G^2$  –  $E^2 > 0$ , then there are two positive solutions. We have

$$
\omega_0 D \sin \omega_0 \tau - \omega_0 F \sin \omega_0 \tau + E \cos \omega_0 \tau
$$
  
+  $FG \cos \omega_0 \tau = \omega_0^2 + mG$   

$$
\omega_0 D \cos \omega_0 \tau - \omega_0 F \cos \omega_0 \tau - FG \sin \omega_0 \tau
$$
  
-  $E \sin \omega_0 \tau = \omega_0 G - \omega_0 m$  (84)

So

$$
\sin \omega_0 \tau = \frac{\left(\omega_0^2 + mG\right)\left(\omega_0 D - \omega_0 F\right) - \omega_0 (G - m)}{\left(\omega_0 D - \omega_0 F\right)^2 - (FG + E)^2}
$$
  

$$
\cos \omega_0 \tau = \frac{-\left(\omega_0^2 + mG\right)(FG + E) + \omega_0 (G - m)}{\left(\omega_0 D - \omega_0 F\right)^2 - (FG + E)^2}
$$
(85)

Then, we get

$$
\tan \omega_0 \tau = \frac{\left(\omega_0^2 + mG\right)\left(\omega_0 D - \omega_0 F\right) - \omega_0 (G - m)}{-\left(\omega_0^2 - mG\right)(FG + E) + \omega_0 (G - m)}.
$$
 (86)

It shows that if  $-2EFH - F^2G^2 + m^2G^2 - E^2 < 0$ ,

$$
\lambda^{2} - (G - ay_{*}e^{-\lambda \tau_{1}} + ca x_{*}e^{-\lambda \tau_{2}} - m)\lambda - mG
$$
  
+ 
$$
may_{*}e^{-\lambda \tau_{1}} + ca x_{*}e^{-\lambda \tau_{2}}G = 0
$$
 (87)

has a pair of imaginary eigenvalues  $\lambda = \pm i\omega_0$ ,  $\eta(\overline{\tau}_1)=0$ 

when  $\overline{\tau}_{j}^{\pm} = (1/\omega_0^{\pm}) \arctan(((\omega_0^2 + mG)(\omega_0^{\pm}D - {\omega_0}^{\pm}F) \omega_0^{\pm}(G - m) / (-(\omega_0^2 + mG)(FG + E) + \omega_0^{\pm}(G - m))) +$  $\pi j/\omega_0^{\pm}, \ \ j=0,1,2,\cdots.$ 

Next verify the cross-sectional conditions:

$$
\left(\frac{d\text{Re}\left(\lambda\right)}{d\tau}\right)^{-1} \neq 0\tag{88}
$$

According to

$$
\lambda^{2} - (G - ay_{*}e^{-\lambda \tau_{1}} + cax_{*}e^{-\lambda \tau_{2}} - m)\lambda - mG
$$
  
+
$$
may_{*}e^{-\lambda \tau_{1}} + cax_{*}e^{-\lambda \tau_{2}}G = 0
$$
 (89)

At this time,  $\tau = \overline{\tau}$ , where  $\overline{\tau}$  is the value of  $\overline{\tau}$  at  $j = 0$ . We get

$$
2\lambda \frac{d\lambda}{d\tau} - \frac{d\lambda}{d\tau} \left( G - De^{-\lambda \tau} + Fe^{-\lambda \tau} - m \right)
$$
  
+ 
$$
\lambda De^{-\lambda \tau} \left( -\tau \frac{d\lambda}{d\tau} - \lambda \right) \lambda
$$
  
- 
$$
\lambda Fe^{-\lambda \tau} \left( -\tau \frac{d\lambda}{d\tau} - \lambda \right) + E e^{-\lambda \tau} \left( -\tau \frac{d\lambda}{d\tau} - \lambda \right)
$$
  
+ 
$$
FG e^{-\lambda \tau} \left( -\tau \frac{d\lambda}{d\tau} - \lambda \right) = 0
$$
 (90)

Then

$$
\left(\frac{d\lambda}{d\tau}\right)^{-1} = -\frac{\tau}{\lambda} - \frac{F - D}{\lambda(\lambda D - \lambda F + E + FG)}
$$

$$
+ \frac{2\lambda - (G - m)}{\lambda[\lambda^2 - \lambda(G - m) - mG]}
$$

$$
\left(\frac{d\text{Re}(\lambda)}{d\tau}\right)^{-1} = \frac{-(F - D)^2}{\omega_0^2(F - D)^2 + (FG + E)^2}
$$
(91)
$$
+ \frac{2(-mG - \omega_0^2) - (G - m)^2}{\omega_0^2(G - m)^2 + (-mG - \omega_0^2)^2}
$$

$$
\neq 0.
$$

## <span id="page-8-0"></span>**4. Numerical Simulations**

In this section, we present some numerical simulations to illustrate our theoretical analysis.

First, we theoretically analyze a predator-prey model with Allee efect in the article and obtain the stability conditions of  $E_1 = (K, 0)$  and  $E_* = (x_*, y_*)$ . Secondly, we carry out numerical simulation and select appropriate the parameters, which draw the stable positions of the equilibrium points  $E_1 = (K, 0)$  and  $E_* = (x_*, y_*)$ , shown in Figures [1](#page-9-0) and [2.](#page-9-1) In the analysis later in this chapter, we use  $A$  as the bifurcation parameter to obtain the critical value of the Hopf bifurcation generated by the model [\(2\).](#page-0-0) By comparing Figures [3](#page-9-2) and [4,](#page-9-3) we find that as the parameter  $A$  changes the equilibrium point changes from a steady state to a limit cycle. To further verify our point of view, we have made a bifurcation diagram as shown in Figure [5.](#page-10-0) We found that the bifurcation parameter produced a bifurcation at about 100, which coincided with our previous guess. By the same analysis method, we give a set of timing diagrams for comparison as shown in Figures [6](#page-10-1) and [7;](#page-10-2) we found that with the increase of the parameter  $A$  the model [\(2\)](#page-0-0) is shown in Figure [6](#page-10-1) and the population gradually becomes stable with the increase of time, while Figure [7](#page-10-2) is a periodic change.

Next, we performed a numerical simulation of the model [\(3\),](#page-1-1) mainly to study the effect of the time-delay parameter  $\tau$ on the stability of the coexistence equilibrium point. Here, we compare three sets of timing diagrams, which are the effect of  $\tau_1$  on the stability of the coexistence equilibrium point, the influence of  $\tau_2$  on the stability of the coexistence equilibrium point, and the influence of  $\tau_1 = \tau_2 = \tau$  on the stability of the coexistence equilibrium point. First, by comparing Figures [8](#page-10-3) and [9,](#page-11-0) we find that when  $\tau_1 = 0.1$ , the population gradually becomes stable with the increase of time. When  $\tau_1$  = 0.25, the population gradually shows periodicity with the increase of time. By comparing Figures [10](#page-11-1) and [11,](#page-11-2) we find that when  $\tau_2 = 0.1$ , the population gradually becomes stable with the increase of time. When  $\tau_2 = 1.5$ , the



<span id="page-9-0"></span>FIGURE 1: The phase portrait of model [\(2\)](#page-0-0) with weak Allee effect. The parameters are taken as  $A = 0.01$ ,  $r = 2.65$ ,  $K = 900$ ,  $a = 0.002$ ,  $c = 0.215$ , and  $m = 1.06$ ;  $E_1 = (K, 0)$  is locally asymptotically stable.



<span id="page-9-1"></span>FIGURE 2: The phase portrait of model [\(2\)](#page-0-0) with weak Allee effect. The parameters are taken as  $A = 0.01$ ,  $r = 2.65$ ,  $K = 900$ ,  $a = 0.02$ ,  $c = 0.215$ , and  $m = 1.06$ ;  $E_* = (x_*, y_*)$  is locally asymptotically stable.

population gradually shows periodicity with the increase of time; by comparing Figures [12](#page-11-3) and [13,](#page-12-0) we found that when  $\tau = 0.1$ , the population gradually becomes stable with the increase of time. When  $\tau = 0.19$ , the population gradually shows periodicity with the increase of time. At the same time, we selected the appropriate parameters and gave three sets of phase diagrams for comparison. By comparing Figures [14](#page-12-1) and [15,](#page-12-2) we find that when  $\tau_1$  increases from 0.1 to 0.25, the coexistence equilibrium point changes from a steady state to a limit cycle; by comparing Figures [16](#page-12-3) and [17,](#page-13-7) we fnd that  $\tau_2$  is increased from 0.1 to 1.5. At this time, the coexistence equilibrium point changes from a steady state to a limit cycle; by comparing Figures [18](#page-13-8) and [19,](#page-13-9) we find that when  $\tau$  increases



<span id="page-9-2"></span>FIGURE 3: The phase portrait of model [\(2\)](#page-0-0) with weak Allee effect. The parameters are taken as  $x(0) = 260$ ,  $y(0) = 70$ ,  $A = 0.01$ ,  $r = 2.65$ ,  $K = 900$ ,  $a = 0.02$ ,  $c = 0.215$ , and  $m = 1.06$ ;  $E_* = (x_*, y_*)$  is locally asymptotically stable, with no limit cycle.



<span id="page-9-3"></span>FIGURE 4: The phase portrait of mode[l \(2\)](#page-0-0) with weak Allee effect. The parameters are taken as  $x(0) = 260$ ,  $y(0) = 70$ ,  $A = 150$ ,  $r = 2.65$ ,  $K = 900$ ,  $a = 0.02$ ,  $c = 0.215$ , and  $m = 1.06$ ;  $E_* = (x_*, y_*)$  becomes unstable; a limit cycle is formed.

from 0.1 to 0.19, the coexistence equilibrium point changes from a steady state to a limit cycle. From this, we can conclude that the stability of the equilibrium point of the model [\(3\)](#page-1-1) changes with the increase of the time lag when the time-delay parameter is introduced, and the generation of the limit cycle (periodic solution) is accompanied by this change.

## **5. Conclusions**

In this paper, we establish a predator-prey model with a weak Allee efect and demonstrate and analyze the boundedness and stability of the model. We also prove that the model



<span id="page-10-0"></span>FIGURE 5: Bifurcation diagram with respect to the parameter  $A$ ; other parameter values are  $r = 2.65$ ,  $K = 900$ ,  $a = 0.02$ ,  $c = 0.215$ , and  $m = 1.06$ .



<span id="page-10-1"></span>FIGURE 6: A Time series diagram for prey and predator. The figure depicts local stability of the interior equilibrium for the model [\(2\),](#page-0-0) where the parameter values are  $x(0) = 230$ ,  $y(0) = 95$ ,  $A = 0.01$ ,  $r = 2.65, K = 900, a = 0.02, c = 0.215, and m = 1.06.$ 

[\(2\)](#page-0-0) experiences transcritical bifurcation around the axial equilibrium and the model [\(2\)](#page-0-0) undergoes a Hopf bifurcation at  $E_* = (x_*, y_*)$  when  $A = A_0$ ; we also analyzed the direction and stability of Hopf bifurcation. Immediately afer we introduced the searching delay and digestion delay in the model [\(2\),](#page-0-0) a new model was obtained, and the model [\(3\)](#page-1-1) was analyzed for stability changes caused by time lag. It is concluded that the stability of the coexistence equilibrium point of the model [\(3\)](#page-1-1) changes as the time lag increases. Finally, we verify our theoretical derivation by numerical simulation. First, we select the appropriate parameters to

satisfy the stable conditions that we deduced in the article and obtain the stable phase diagrams of the equilibrium points  $E_1 = (K, 0)$  and  $E_* = (x_*, y_*)$ . Second, we try to change the value of parameter  $A$ , the timing diagram, phase diagram, and bifurcation diagram corresponding to each parameter drawn. Further verifcation of our conclusion is as  $A$  increases, the model  $(2)$  will produce bifurcation. We

<span id="page-10-3"></span>depicts local stability of the interior equilibrium for the delayed model [\(3\)](#page-1-1) with the time delays  $\tau_1 = 0.1$  and  $\tau_2 = 0$ , where the other parameter values are  $x(0) = 230$ ,  $y(0) = 95$ ,  $A = 0.01$ ,  $r = 2.65$ ,

 $K = 900$ ,  $a = 0.02$ ,  $c = 0.215$ , and  $m = 1.06$ .



<span id="page-10-2"></span>Figure 7: A Time series diagram for prey and predator. Existence of periodic solution around the interior equilibrium  $E_* = (x_*, y_*)$  for the model [\(2\),](#page-0-0) where the parameter values are  $x(0) = 230, y(0) = 95$ ,  $A = 150, r = 2.65, K = 900, a = 0.02, c = 0.215, and m = 1.06.$ 





<span id="page-11-0"></span>Figure 9: A Time series diagram for prey and predator. Existence of periodic solution around the interior equilibrium  $E_* = (x_*, y_*)$  for delayed model [\(3\)](#page-1-1) with the time delays  $\tau_1 = 0.25$  and  $\tau_2 = 0$ , where the parameter values are  $x(0) = 230$ ,  $y(0) = 95$ ,  $A = 0.01$ ,  $r = 2.65$ ,  $K = 900$ ,  $a = 0.02$ ,  $c = 0.215$ , and  $m = 1.06$ .





<span id="page-11-2"></span>Figure 11: A Time series diagram for prey and predator. Existence of periodic solution around the interior equilibrium  $E_* = (x_*, y_*)$  for delayed model [\(3\)](#page-1-1) with the time delays  $\tau_1 = 0$  and  $\tau_2 = 1.5$ , where the parameter values are  $x(0) = 230$ ,  $y(0) = 95$ ,  $A = 0.01$ ,  $r = 2.65$ ,  $K = 900$ ,  $a = 0.02$ ,  $c = 0.215$ , and  $m = 1.06$ .





<span id="page-11-1"></span>FIGURE 10: A Time series diagram for prey and predator. The figure depicts local stability of the interior equilibrium for the delayed model [\(3\)](#page-1-1) with the time delays  $\tau_1 = 0$  and  $\tau_2 = 0.1$ , where the other parameter values are  $x(0) = 230$ ,  $y(0) = 95$ ,  $A = 0.01$ ,  $r = 2.65$ ,  $K = 900$ ,  $a = 0.02$ ,  $c = 0.215$ , and  $m = 1.06$ .

<span id="page-11-3"></span>FIGURE 12: A Time series diagram for prey and predator. The figure depicts local stability of the interior equilibrium for the delayed model [\(3\)](#page-1-1) with the time delay  $\tau_1 = \tau_2 = 0.1$ , where the other parameter values are  $x(0) = 230$ ,  $y(0) = 95$ ,  $A = 0.01$ ,  $r = 2.65$ ,  $K = 900$ ,  $a = 0.02$ ,  $c = 0.215$ , and  $m = 1.06$ .

also carry out a numerical model of the model [\(3\)](#page-1-1) and discuss our numerical results in three groups: at the beginning, let  $\tau_1$ change,  $\tau_2 = 0$ ; we observe the timing diagram corresponding to the model [\(3\)](#page-1-1) as  $\tau_1$  increases. In the change of the phase diagram, we find that with the increase of  $\tau_1$  the coexistence equilibrium point of the model [\(3\)](#page-1-1) begins to stabilize and

becomes unstable and is also accompanied by the generation of the limit cycle (periodic solution); we also change the  $\tau_2$ ,  $\tau_1$  = 0; we observe the increase of  $\tau_2$ , corresponding to the change of the phase diagram of the time series of the model [\(3\);](#page-1-1) we find that with the increase of  $\tau_2$  the coexistence equilibrium point of the model [\(3\)](#page-1-1) begins to stabilize and



<span id="page-12-0"></span>Figure 13: A Time series diagram for prey and predator. Existence of periodic solution around the interior equilibrium  $E_* = (x_*, y_*)$ for delayed model [\(3\)](#page-1-1) with the time delay  $\tau_1 = \tau_2 = 0.19$ , where the parameter values are  $x(0) = 230$ ,  $y(0) = 95$ ,  $A = 0.01$ ,  $r = 2.65$ ,  $K = 900$ ,  $a = 0.02$ ,  $c = 0.215$ , and  $m = 1.06$ .



<span id="page-12-1"></span>FIGURE 14: The phase portrait of delayed model [\(3\)](#page-1-1) with the time delays  $\tau_1 = 0.1$  and  $\tau_2 = 0$ . The parameters are taken as  $x(0) = 230$ ,  $y(0) = 95$ ,  $A = 0.01$ ,  $r = 2.65$ ,  $K = 900$ ,  $a = 0.02$ ,  $c = 0.215$ , and  $m = 1.06; E_* = (x_*, y_*)$  is locally asymptotically stable, with no limit cycle.

becomes unstable. At the same time, it is accompanied by the generation of the limit cycle (periodic solution); the same we let  $\tau_1 = \tau_2 = \tau$ , so that  $\tau$  increases; we also found similar changes. In this process, we also found an interesting phenomenon. When  $\tau_1$  takes a small value, the stability of the coexistence equilibrium point of the model changes. However,  $\tau_2$  requires a larger value than  $\tau_1$ . When  $\tau_1$  =  $\tau_2$  =  $\tau$ ,  $\tau$  only needs to take a small value, and the stability of the coexistence equilibrium point of the model

changes. From a biological point of view, we fnd that the introduction of weak Allee effect will change the stability of the model; that is to say, the stable state between populations will be broken. Similarly, delays can destroy the stability of the original predator-prey model. Moreover, the introduction of delay and weak Allee efect makes the model closer to reality and makes us more accurately understand the dynamic changes of interspecific relationships. Therefore, we can find that Allee efect and delay play an important role in describing population dynamics.



<span id="page-12-2"></span>FIGURE 15: The phase portrait of delayed model [\(3\)](#page-1-1) with the time delays  $\tau_1 = 0.25$  and  $\tau_2 = 0$ . The parameters are taken as  $x(0) = 230$ ,  $y(0) = 95$ ,  $A = 0.01$ ,  $r = 2.65$ ,  $K = 900$ ,  $a = 0.02$ ,  $c = 0.215$ , and  $m = 1.06; E_* = (x_*, y_*)$  becomes unstable; a limit cycle is formed.



<span id="page-12-3"></span>FIGURE 16: The phase portrait of delayed model [\(3\)](#page-1-1) with the time delays  $\tau_1 = 0$  and  $\tau_2 = 0.1$ . The parameters are taken as  $x(0) =$  $230, y(0) = 95, A = 0.01, r = 2.65, K = 900, a = 0.02, c = 0.02$ 0.215, and  $m = 1.06$ ;  $E_* = (x_*, y_*)$  is locally asymptotically stable, with no limit cycle.



<span id="page-13-7"></span>FIGURE 17: The phase portrait of delayed model [\(3\)](#page-1-1) with the time delays  $\tau_1 = 0$  and  $\tau_2 = 1.5$ . The parameters are taken as  $x(0) = 230$ ,  $y(0) = 95$ ,  $A = 0.01$ ,  $r = 2.65$ ,  $K = 900$ ,  $a = 0.02$ ,  $c = 0.215$ , and  $m = 1.06; E_* = (x_*, y_*)$  becomes unstable; a limit cycle is formed.



<span id="page-13-8"></span>FIGURE 18: The phase portrait of delayed model [\(3\)](#page-1-1) with the time delay  $\tau_1 = \tau_2 = 0.1$ . The parameters are taken as  $x(0) = 230$ ,  $y(0) = 0.1$ 95,  $A = 0.01$ ,  $r = 2.65$ ,  $K = 900$ ,  $a = 0.02$ ,  $c = 0.215$ , and  $m = 1.06$ ;  $E_* = (x_*, y_*)$  is locally asymptotically stable, with no limit cycle.

## **Data Availability**

The data used to support the findings of this study are included within the article.

## **Conflicts of Interest**

The authors declare that they have no conflicts of interest.

## **Acknowledgments**

This work was supported by the National Natural Science Foundation of China (31260098, 31560127), the Fundamental Research Funds for the Central Universities (31920180116, 31920180044, and 31920170072), the Program for Yong Talent



<span id="page-13-9"></span>FIGURE 19: The phase portrait of delayed model [\(3\)](#page-1-1) with the time delay  $\tau_1 = \tau_2 = 0.19$ . The parameters are taken as  $x(0) = 230$ ,  $y(0) = 0.19$ . 95,  $A = 0.01$ ,  $r = 2.65$ ,  $K = 900$ ,  $a = 0.02$ ,  $c = 0.215$ , and  $m = 1.06$ ;  $E_* = (x_*, y_*)$  becomes unstable; a limit cycle is formed.

of State Ethnic Afairs Commission of China (No. [2014]121), Gansu Provincial First-Class Discipline Program of Northwest Minzu University, and Central Universities Fundamental Research Funds for the Graduate Students of Northwest Minzu University (Yxm2019109).

## **References**

- <span id="page-13-0"></span>[1] V. Volterra, "Fluctuations in the abundance of a species considered mathematically," *Nature*, vol. 118, no. 2972, pp. 558–560, 1926.
- <span id="page-13-1"></span>[2] A. J. Lotka, *Elements of Physical Biology*, Williams and Wilkins, Pennsylvania, Pa, USA, 1926.
- <span id="page-13-2"></span>[3] F. Courchamp, T. Clutton-Brock, and B. Grenfell, "Inverse density dependence and the Allee efect," *Trends in Ecology & Evolution*, vol. 14, no. 10, pp. 405–410, 1999.
- [4] P. A. Stephens, W. J. Sutherland, and R. P. Freckleton, "What is the allee efect?" *Oikos*, vol. 87, no. 1, pp. 185–190, 1999.
- <span id="page-13-3"></span>[5] P. A. Stephens and W. J. Sutherland, "Consequences of the allee efect for behaviour, ecology and conservation," *Trends in Ecology & Evolution*, vol. 14, no. 10, pp. 401–405, 1999.
- <span id="page-13-4"></span>[6] W. C. Allee, *Animal Aggregations: A study in General Sociology*, University of Chicago Press, Illinois, Ill, USA, 1931.
- <span id="page-13-5"></span>[7] A. Deredec and F. Courchamp, "Extinction thresholds in hostparasite dynamics," *Annales Zoologici Fennici*, vol. 40, no. 2, pp. 115–130, 2003.
- <span id="page-13-6"></span>[8] M.-H. Wang and M. Kot, "Speeds of invasion in a model with strong or weak Allee efects," *Mathematical Biosciences*, vol. 171, no. 1, pp. 83–97, 2001.
- [9] C. W. Fowler and J. D. Baker, "A review of animal population dynamics at extremely reduced population levels," *Report - International Whaling Commission*, vol. 41, pp. 545–554, 1991.
- [10] S. J. Schreiber, "Allee effects, extinctions, and chaotic transients in simple population models," *Theoretical Population Biology*, vol. 64, no. 2, pp. 201–209, 2003.
- [11] M. Wang, M. Kot, and M. G. Neubert, "Integrodiference equations, Allee efects, and invasions," *Journal of Mathematical Biology*, vol. 44, no. 2, pp. 150–168, 2002.
- <span id="page-14-5"></span>[12] G. Wang, X.-G. Liang, and F.-Z. Wang, "The competitive dynamics of populations subject to an Allee efect," *Ecological Modelling*, vol. 124, no. 2-3, pp. 183–192, 1999.
- <span id="page-14-0"></span>[13] R. Lin, S. Liu, and X. Lai, "Bifurcations of a predator-prey system with weak allee efects," *Journal of the Korean Mathematical Society*, vol. 50, no. 4, pp. 695–713, 2013.
- <span id="page-14-1"></span>[14] S.-R. Zhou, Y.-F. Liu, and G. Wang, "The stability of predatorprey systems subject to the Allee effects," *Theoretical Population Biology*, vol. 67, no. 1, pp. 23–31, 2005.
- [15] L. Shi, H. Liu, Y. Wei, M. Ma, and J. Ye, "The permanence and periodic solution of a competitive system with infnite delay, feedback control, and Allee efect," *Advances in Difference Equations*, vol. 2018, no. 1, article 400, 2018.
- <span id="page-14-2"></span>[16] H. Liu, Z. Li, M. Gao, H. Dai, and Z. Liu, "Dynamics of a hostparasitoid model with Allee effect for the host and parasitoid aggregation," *Ecological Complexity*, vol. 6, no. 3, pp. 337–345, 2009.
- <span id="page-14-3"></span>[17] J. Wang, J. Shi, and J. Wei, "Predator-prey system with strong Allee efect in prey," *Journal of Mathematical Biology*, vol. 62, no. 3, pp. 291–331, 2011.
- [18] J. Wang, J. Shi, and J. Wei, "Dynamics and pattern formation in a difusive predator-prey system with strong Allee efect in prey," *Journal of Differential Equations*, vol. 251, no. 4-5, pp. 1276–1304, 2011.
- <span id="page-14-4"></span>[19] Y. Cai, C. Zhao, W. Wang, and J. Wang, "Dynamics of a Leslie-Gower predator-prey model with additive Allee effect," *Applied Mathematical Modelling: Simulation and Computation for Engineering and Environmental Systems*, vol. 39, no. 7, pp. 2092–2106, 2015.
- <span id="page-14-6"></span>[20] C. Li and H. Zhu, "Canard cycles for predator-prey systems with holling types of functional response," *Journal of Differential Equations*, vol. 254, no. 2, pp. 879–910, 2013.
- [21] Y. Li and D. Xiao, "Bifurcations of a predator-prey system of holling and leslie types," *Chaos, Solitons & Fractals*, vol. 34, no. 2, pp. 606–620, 2007.
- [22] Y. Qu and J. Wei, "Bifurcation analysis in a time-delay model for prey-predator growth with stage-structure," *Nonlinear Dynamics*, vol. 49, no. 1-2, pp. 285–294, 2007.
- [23] J. P. Tripathi, S. Tyagi, and S. Abbas, "Global analysis of a delayed density dependent predator-prey model with Crowley-Martin functional response," *Communications in Nonlinear Science and Numerical Simulation*, vol. 30, no. 1–3, pp. 45–69, 2016.
- [24] J. Ylikarjula, S. Alaja, J. Laakso, and D. Tesar, "Effects of patch number and dispersal patterns on population dynamics and synchrony," *Journal of Theoretical Biology*, vol. 207, no. 2-3, pp. 377–387, 2000.
- [25] X. Wang and J. Wei, "Dynamics in a difusive predatorprey system with strong Allee effect and Ivlev-type functional response," *Journal of Mathematical Analysis and Applications* , vol. 422, no. 2, pp. 1447–1462, 2015.
- <span id="page-14-7"></span>[26] Z. Ma, H. Tang, S. Wang, and T. Wang, "Bifurcation of a predator-prey system with generation delay and habitat complexity," *Journal of the Korean Mathematical Society*, vol. 55, no. 1, pp. 43–58, 2018.



International Journal of [Mathematics and](https://www.hindawi.com/journals/ijmms/)  **Mathematical Sciences** 

ww.hindawi.com Volume 2018 / Mary 2018





[Applied Mathematics](https://www.hindawi.com/journals/jam/)

www.hindawi.com Volume 2018



**The Scientifc [World Journal](https://www.hindawi.com/journals/tswj/)**



[Probability and Statistics](https://www.hindawi.com/journals/jps/) Hindawi www.hindawi.com Volume 2018 Journal of







Engineering [Mathematics](https://www.hindawi.com/journals/ijem/)

International Journal of

[Complex Analysis](https://www.hindawi.com/journals/jca/) www.hindawi.com Volume 2018

www.hindawi.com Volume 2018 [Stochastic Analysis](https://www.hindawi.com/journals/ijsa/) International Journal of



www.hindawi.com Volume 2018 Advances in<br>[Numerical Analysis](https://www.hindawi.com/journals/ana/)



www.hindawi.com Volume 2018 **[Mathematics](https://www.hindawi.com/journals/jmath/)** 



[Submit your manuscripts at](https://www.hindawi.com/) www.hindawi.com

Hindawi

 $\bigcirc$ 

www.hindawi.com Volume 2018 [Mathematical Problems](https://www.hindawi.com/journals/mpe/)  in Engineering Advances in **Discrete Dynamics in** Mathematical Problems and International Journal of **Discrete Dynamics in** 



Journal of www.hindawi.com Volume 2018 [Function Spaces](https://www.hindawi.com/journals/jfs/)



Differential Equations International Journal of



Abstract and [Applied Analysis](https://www.hindawi.com/journals/aaa/) www.hindawi.com Volume 2018



Nature and Society



www.hindawi.com Volume 2018 <sup>Advances in</sup><br>[Mathematical Physics](https://www.hindawi.com/journals/amp/)